



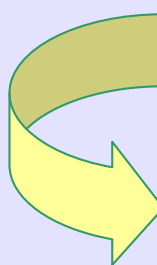
Beauty Hadrons:

- **Lifetime Ratios:** $\tau(\mathbf{B}^+)/\tau(\mathbf{B}_d), \tau(\mathbf{B}_s)/\tau(\mathbf{B}_d), \tau(\Lambda_b)/\tau(\mathbf{B}_d)$
- **Width Differences:** $\Delta\Gamma_d, \Delta\Gamma_s$
- **Semileptonic CP-Asymmetries:** A_{SL}^d, A_{SL}^s

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$\mathbf{B}_q^0 - \overline{\mathbf{B}}_q^0$ **SYSTEMS** ($q=d,s$)

$$\hat{H} = \hat{M} - \frac{i}{2} \hat{\Gamma}$$



$$\hat{M} = \begin{pmatrix} M_{11} & M_{12} \\ M_{12}^* & M_{11} \end{pmatrix}$$

Oscillation



$$\hat{\Gamma} = \begin{pmatrix} \Gamma_{11} & \Gamma_{12} \\ \Gamma_{12}^* & \Gamma_{11} \end{pmatrix}$$

Decay

Eigenstates:

$$|\mathbf{B}_q^{L,H}\rangle = p_q |\mathbf{B}_q^0\rangle \pm q_q |\overline{\mathbf{B}}_q^0\rangle$$

$$\Gamma_{11}^q \quad (\Delta B = 0)$$

• Lifetimes: $\Gamma_{B_q} = \Gamma_{11}^q$

$$\Gamma_{12}^q \quad (\Delta B = 2)$$

• Width Differences:

$$\Delta\Gamma_q = \Gamma_{B_L^q} - \Gamma_{B_H^q} \leftrightarrow \Gamma_{12}^q$$

different

• Semileptonic CP-Asymmetries:

$$\mathbf{A}_{SL}^q = -2 \left(\left| \frac{q}{p} \right|_q - 1 \right) \leftrightarrow \Gamma_{12}^q$$

V_{CKM}

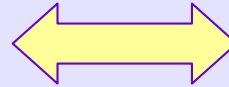
$$\Gamma_{11}^q \propto \text{Disc} \langle \mathbf{B}_q^0 | T | \mathbf{B}_q^0 \rangle$$

$$\Gamma_{12}^q \propto \text{Disc} \langle \mathbf{B}_q^0 | T | \overline{\mathbf{B}}_q^0 \rangle$$

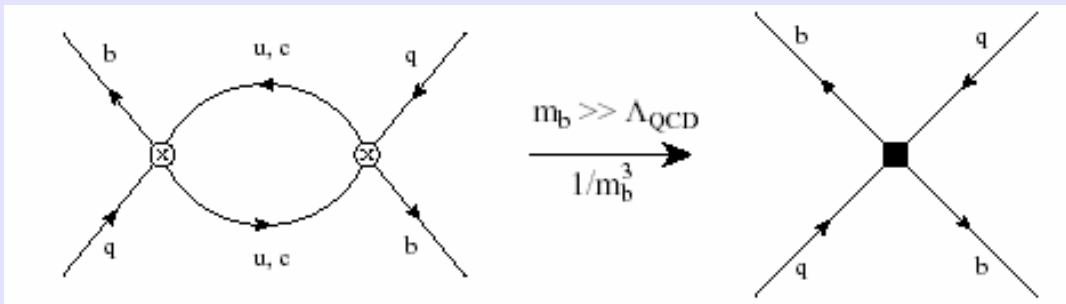
$$T = i \int d^4x \mathcal{T}(H_{\text{eff}}^{\Delta B=1}(\mathbf{x}) H_{\text{eff}}^{\Delta B=1}(0))$$

HQE ($m_b \gg \Lambda_{\text{QCD}}$)

Large energy release



contact interaction



SPECTATOR EFFECTS

$$\Gamma_{12}^q = \sum_k \frac{\vec{c}_k^q(\mu)}{m_b^k} \langle \mathbf{B}_q^0 | \vec{O}_k^{q\Delta B=2}(\mu) | \overline{\mathbf{B}}_q^0 \rangle$$

Scale separation

- $\vec{c}_k^q(\mu)$: short-distance (perturbative)
- $\langle \mathbf{B}_q^0 | \vec{O}_k^{q\Delta B=2}(\mu) | \overline{\mathbf{B}}_q^0 \rangle$: long-distance (non-perturbative)

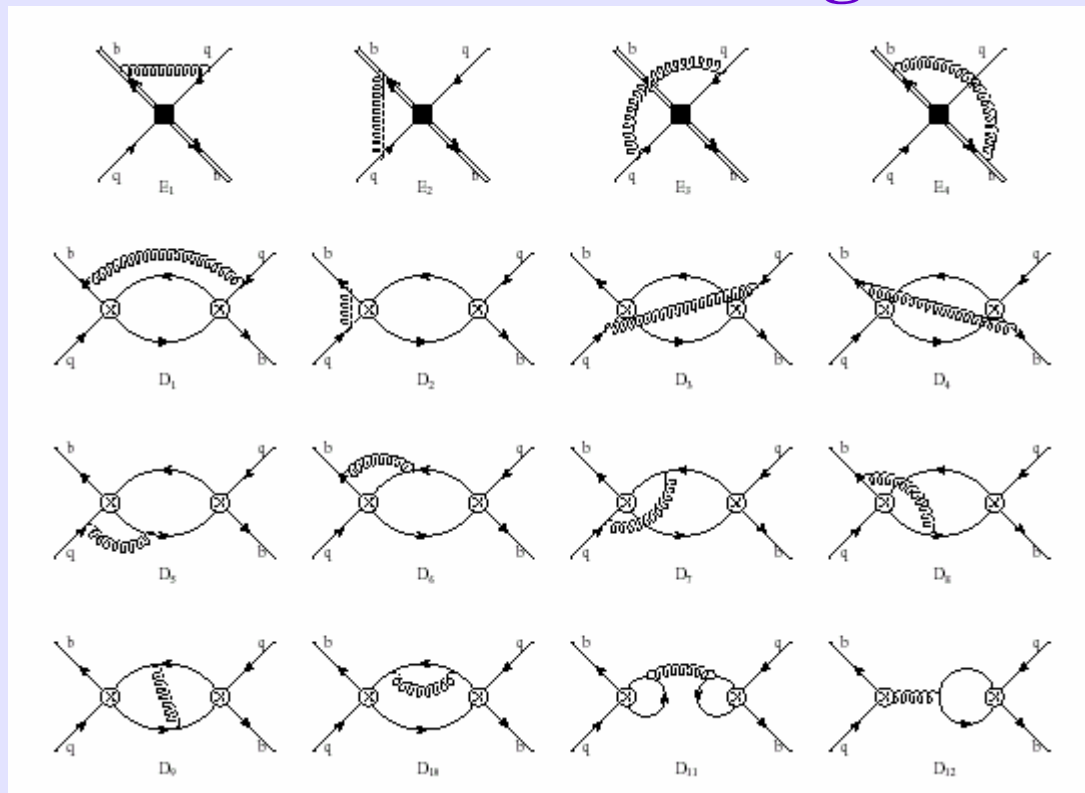
Wilson coefficients

$$\text{Disc}\langle T \rangle = \vec{c}(\mu) \langle \vec{O}(\mu) \rangle$$

Matching at NLO in QCD and $O(1/m_b)$ in the HQE

$$\Delta B = 2(\Delta B = 0)$$

NLO Diagrams



• Lifetime Ratios

$$\Gamma_{11}^q = \frac{G_F^2 |V_{cb}|^2 m_b^5}{192\pi^3 (2M_{B_q})} \left[c^{(3)} \langle \bar{b}b \rangle + c^{(5)} \frac{g_s}{m_b^2} \langle \bar{b} \sigma_{\mu\nu} G^{\mu\nu} b \rangle + \frac{96\pi^2}{m_b^3} \sum_k \left(c_k^{(6)} \langle O_k^{(6)} \rangle + \frac{c_k^{(7)}}{m_b} \langle O_k^{(7)} \rangle \right) \right]$$

$O(1)$ (1996) [M. Neubert and C.T. Sachrajda]

$O(\alpha_s)$ (2002)

• [E. Franco, V. Lubicz, F. Mescia and C.T.]

• [M. Beneke, G. Buchalla, C. Greub, A. Lenz and U. Nierste]

**Spectator effects,
explaining lifetime diversity,
appear at $O(1/m_b^3)$**

$O(1/m_b)$ (2004)

• [F. Gabbiani, A. I. Onishchenko and A. A. Petrov]

• Width Differences and \mathcal{CP}

$$\Gamma_{12}^s = -\frac{G_F^2 m_b^2}{12\pi (2M_{B_s})} (V_{cs}^* V_{cb})^2 \left[C_1 \langle O_1^s \rangle + C_2 \langle O_2^s \rangle + \delta_{1/m_b} \right]$$

$O(\alpha_s)$ (2003)

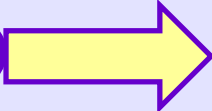
• [M. Ciuchini, E. Franco, V. Lubicz, F. Mescia and C.T.]

• [M. Beneke, G. Buchalla, A. Lenz and U. Nierste]

**The leading contribution
in the HQE is of $O(1/m_b^3)$**

$O(1/m_b)$ (1996)

• [M. Beneke, G. Buchalla and I. Dunietz]

$O(1/m_b^2)$ (coming soon?)  + convenient change of basis

• [A. Lenz and U. Nierste]

[appeared today: hep-ph/0612167]

Matrix Elements: $\Delta B=0$ Operators

Leading spectator effect contribution: $O(1/m_b^3)$

$$O_1^q = (\bar{b}q)_{V-A} (\bar{q}b)_{V-A} \leftrightarrow B_1^q,$$

$$O_2^q = (\bar{b}q)_{S-P} (\bar{q}b)_{S+P} \leftrightarrow B_2^q,$$

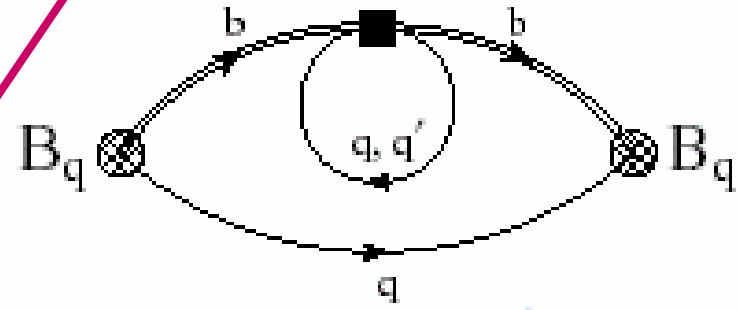
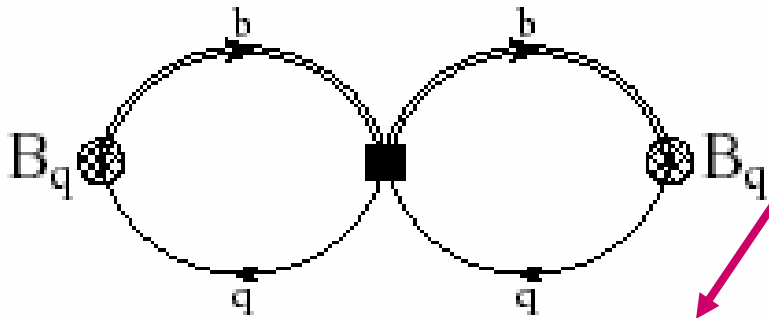
$$O_3^q = (\bar{b}T^a q)_{V-A} (\bar{q}T^a b)_{V-A} \leftrightarrow \epsilon_1^q,$$

$$O_4^q = (\bar{b}T^a q)_{S-P} (\bar{q}T^a b)_{S+P} \leftrightarrow \epsilon_2^q,$$

$$O_P = (\bar{b}T^a b)_V \sum_{q=u,d,s,c} (\bar{q}T^a q)_V$$

not computed

$$[(\bar{q}q)_{V-A} = \bar{q}\gamma_L^\mu q, (\bar{q}q)_{S\pm P} = \bar{q}(1 \pm \gamma_5)q, (\bar{q}q)_V = \bar{q}\gamma^\mu q]$$



These contributions are expected to be:

- negligible in $\tau(B^+)/\tau(B_d)$ [~~SU(2)~~-effect]
- small in $\tau(B_s)/\tau(B_d)$ [~~SU(3)~~-effect]
- significant in $\tau(\Lambda_b)/\tau(B_d)$

not computed

B-parameters:

Λ_b Baryon

- **Only one lattice calculation, in HQET ($m_b \rightarrow \infty$), evolved at LO**

[M. Di Pierro, C.T. Sachrajda, C. Michael (UKQCD Collaboration), **1999**]

(B_d, B_s, B^+) Mesons

- **Lattice-HQET ($m_b \rightarrow \infty$)**

$$B_1^d = 1.06 \pm 0.08, \quad B_2^d = 1.01 \pm 0.07,$$

[M. Di Pierro and C.T. Sachrajda, **1998**]

$$\epsilon_1^d = -0.01 \pm 0.03, \quad \epsilon_2^d = -0.03 \pm 0.02.$$

- **Lattice-QCD ($m_c \leq m_Q < m_b, m_Q \rightarrow m_b$)**

$$B_1^d = 1.2 \pm 0.2, \quad B_2^d = 0.9 \pm 0.1,$$

[APE (D. Becirevic et al.), **2001**]

$$\epsilon_1^d = 0.04 \pm 0.01, \quad \epsilon_2^d = 0.04 \pm 0.01.$$

- **Sum Rules, in HQET**

$$B_1^d = 1.01 \pm 0.01, \quad B_2^d = 0.99 \pm 0.01,$$

[M.S. Baek et al., **1998**]

$$\epsilon_1^d = -0.08 \pm 0.02, \quad \epsilon_2^d = -0.01 \pm 0.03.$$

Subleading spectator effect contribution: $O(1/m_b^4)$

8 operators, from the **VSA** (B-mesons) or the **quark-diquark model** (baryon)

Matrix Elements: $\Delta B=2$ Operators

Leading contribution: $O(1/m_b^3)$

$$O_1^q = (\bar{b}q)_{V-A} (\bar{b}q)_{V-A} \leftrightarrow B_1^q,$$

$$O_2^q = (\bar{b}q)_{S-P} (\bar{b}q)_{S-P} \leftrightarrow B_2^q,$$

B-parameters: From QCD *-sum rules* [J.G. Korner et al., 2003] and

Lattice:

• HQET ($m_b \rightarrow \infty$)

$$B_1^s = 0.83(5)(6), \quad B_2^s = 0.81(2)(10)$$

[V. Gimenez and J. Reyes, 2000]

• NRQCD ($O(1/m_b)$)

$$B_1^s = 0.85(3)(11), \quad B_2^s = 0.82(2)(11)$$

[Hi-KEK (S. Hashimoto et al.), 2000]

• Unquenched NRQCD ($n_f=2$)

$$B_1^s = 0.85(2)(6), \quad B_2^s = 0.84(6)(8)$$

[JLQCD (S. Aoki et al.), 2001-2003]

• QCD ($m_c \leq m_Q < m_b, m_Q \rightarrow m_b$)

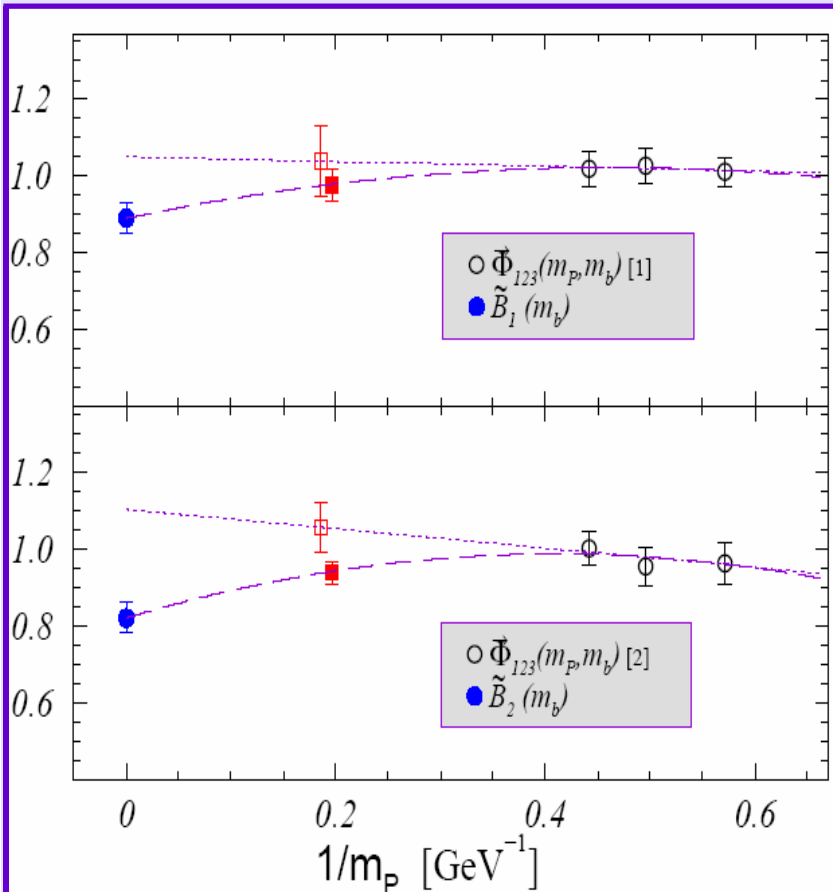
$$B_1^s = 0.91(3)_{-6}^{+0}, \quad B_2^s = 0.86(2)_{-3}^{+2}$$

[APE (D. Becirevic et al.), 2000]

•QCD +HQET

$$\mathbf{B}_1^s = 0.87(2)(5), \quad \mathbf{B}_2^s = 0.84(2)(4)$$

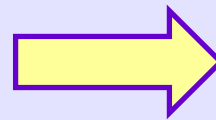
[APE (D. Becirevic et al.), 2001]



Relativistic QCD and static HQET
results are combined

Matching between QCD and HQET
is performed at NLO
in perturbation theory

A new NRQCD study
with $nf=2+1$ (staggered)
finds compatible results



$$\mathbf{B}_1^s = 0.76(11), \quad \mathbf{B}_2^s = 0.84(13)$$

[HPQCD (E.Dalgic et al.), 2006]

Subleading contribution: $O(1/m_b^4)$

4 Operators $\{R_1^q, R_2^q, R_3^q, R_4^q\}$:

- R_1^q, R_4^q are related to operators computed on the lattice, through Fierz Identities and equations of motion
- R_2^q, R_3^q are estimated in the Vacuum Insertion Approximation

A calculation with QCD-sum rules is in progress

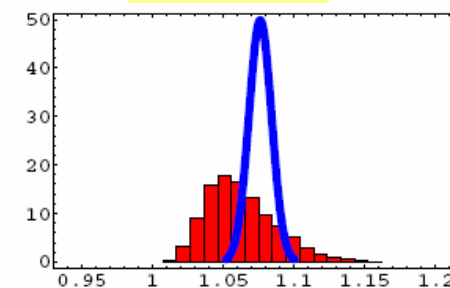
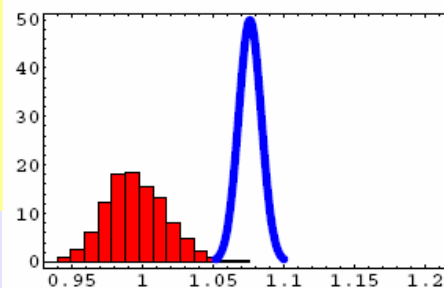
[A.A. Pivovarov, private communication]

Lifetime Ratios

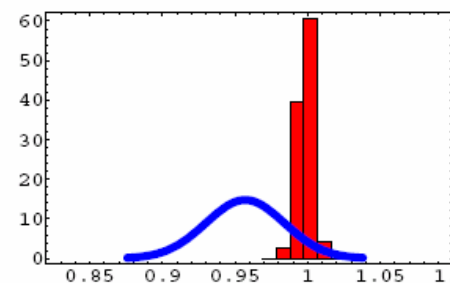
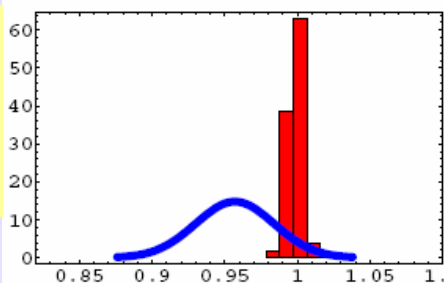
LO

NLO

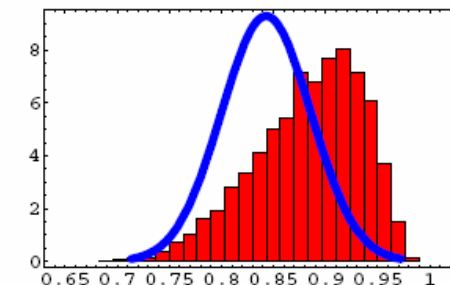
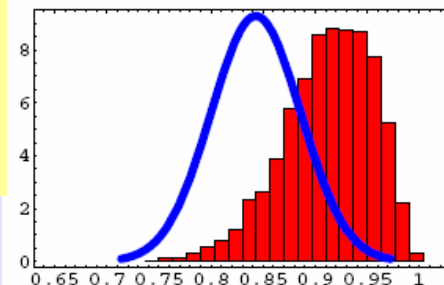
$$\frac{\tau(B^+)}{\tau(B_d)}$$



$$\frac{\tau(B_s)}{\tau(B_d)}$$



$$\frac{\tau(\Lambda_b)}{\tau(B_d)}$$



	$\frac{\tau(B^+)}{\tau(B_d)}$	$\frac{\tau(B_s)}{\tau(B_d)}$	$\frac{\tau(\Lambda_b)}{\tau(B_d)}$
LO	1.01(3)	1.00(1)	0.93(4)
NLO	1.06(3)	1.00(1)	0.90(5)
NLO+ $O(1/m_b^4)$	1.06(2)	1.00(1)	0.88(5)

Theoretical predictions at the NLO + contribution of $O(1/m_b^4)$:

$$\frac{\tau(\mathbf{B}^+)}{\tau(\mathbf{B}_d)} = 1.06 \pm 0.02, \quad \frac{\tau(\mathbf{B}_s)}{\tau(\mathbf{B}_d)} = 1.00 \pm 0.01, \quad \frac{\tau(\Lambda_b)}{\tau(\mathbf{B}_d)} = 0.88 \pm 0.05$$

[E.Franco, V.Lubicz, F.Mescia, C.T., 2002-2004]

Experimental averages:

$$\frac{\tau(\mathbf{B}^+)}{\tau(\mathbf{B}_d)} = 1.076 \pm 0.008, \quad \frac{\tau(\mathbf{B}_s)}{\tau(\mathbf{B}_d)} = 0.957 \pm 0.027, \quad \frac{\tau(\Lambda_b)}{\tau(\mathbf{B}_d)} = 0.84 \pm 0.05$$

[HFAG, 2006]

...before new Tevatron measurements...

CDF

$$\frac{\tau(\Lambda_b)}{\tau(\mathbf{B}_d)} = 1.041 \pm 0.057$$

[CDF (A.Abulencia et al.), 2006]

D0 [PRELIMINARY]

$$\frac{\tau(\Lambda_b)}{\tau(\mathbf{B}_d)} = 0.870 \pm 0.102 \pm 0.041$$

[D0 (C.Ay talk@HQL06), 2006]

- differs from the world average by 2.5σ
- single best measurement from fully reconstructed $\Lambda_b \rightarrow \Lambda^0 J_\psi$
- Is there a reversed Λ_b puzzle? Too early to answer!

Width Differences and Semileptonic CP-Asymmetries

$$\frac{\Delta\Gamma_q}{\Gamma_q} = -\frac{\Delta M_q}{\Gamma_q} \operatorname{Re}\left(\frac{\Gamma_{12}^q}{M_{12}^q}\right)$$

$$A_{\text{SL}}^q = \operatorname{Im}\left(\frac{\Gamma_{12}^q}{M_{12}^q}\right)$$

After the CDF measurement of ΔM_s
non-perturbative uncertainties
 are only enclosed in **B-parameters**

Γ_{12}^q describes **decay** (to SM particles only)

Insensitive to NP

M_{12}^q describes **oscillation**
 through box-diagrams

Sensitive
 to NP

$$M_{12}^q = (M_{12}^q)_{\text{SM}} C_{B_q} e^{2i\varphi_{B_q}}$$

$$\frac{\Delta\Gamma_q}{\Gamma_q} = -\left(\frac{\Delta M_q}{\Gamma_q}\right)^{\text{exp}} \left[\operatorname{Re}\left(\frac{\Gamma_{12}^q}{M_{12}^q}\right)^{\text{SM}} \frac{\cos 2\varphi_{B_q}}{C_{B_q}} - \operatorname{Im}\left(\frac{\Gamma_{12}^q}{M_{12}^q}\right)^{\text{SM}} \frac{\sin 2\varphi_{B_q}}{C_{B_q}} \right]$$

$$A_{\text{SL}}^q = \left[\operatorname{Im}\left(\frac{\Gamma_{12}^q}{M_{12}^q}\right)^{\text{SM}} \frac{\cos 2\varphi_{B_q}}{C_{B_q}} - \operatorname{Re}\left(\frac{\Gamma_{12}^q}{M_{12}^q}\right)^{\text{SM}} \frac{\sin 2\varphi_{B_q}}{C_{B_q}} \right]$$

In particular
 $\operatorname{Re} \gg \operatorname{Im}$ in B_s

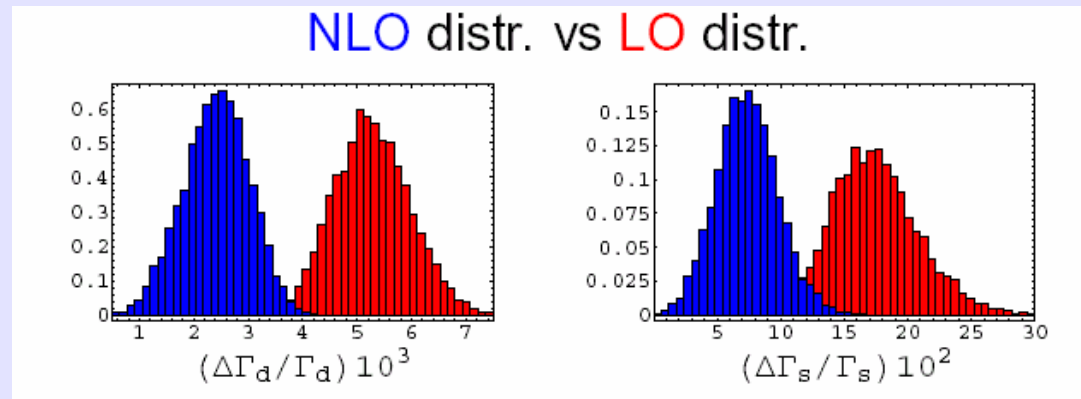
Even a small φ_{B_s}
 can be seen in A_{SL}^s

A promising way to determine $\Delta M_s / (\Delta M_s)_{\text{SM}}$ is provided by:

$$\frac{\Delta M_s}{(\Delta M_s)_{\text{SM}}} = C_{B_s} = -\left| \operatorname{Re}\left(\frac{\Gamma_{12}^s}{M_{12}^s}\right)^{\text{SM}} \right| \frac{S_{\psi\phi}}{A_{\text{SL}}^s}$$

[M. Blanke, A.J. Buras, D. Guadagnoli, C.T., 2006]

Width Differences



Experimental averages:

$$\frac{\Delta\Gamma_d}{\Gamma_d} = (9 \pm 37) \cdot 10^{-3}, \quad \frac{\Delta\Gamma_s}{\Gamma_s} = (14 \pm 6) \cdot 10^{-2}$$

[HFAG, 2006]

[Unofficial world average, talk by
R. Van Kooten@FP&CP2006]

Theoretical predictions at the NLO + contribution of $O(1/m_b^4)$:

$$\frac{\Delta\Gamma_d}{\Gamma_d} = (2.3 \pm 0.8) \cdot 10^{-3}, \quad \frac{\Delta\Gamma_s}{\Gamma_s} = (7 \pm 3) \cdot 10^{-2}$$

[M.Ciuchini, E.Franco, V.Lubicz, F.Mescia, C.T., 2003]

Quite large uncertainties for NLO and $O(1/m_b^4)$ predictions:
they are due to large NLO and $O(1/m_b^4)$ contributions
that induce strong cancellations

A. Lenz's and U. Nierste's Idea: Change Basis!

[Today's hep-ph/0612167]

$$O_1^q = (\bar{b}_i q_i)_{V-A} (\bar{b}_j q_j)_{V-A} \leftrightarrow B_1^q, \quad O_2^q = (\bar{b}_i q_i)_{S-P} (\bar{b}_j q_j)_{S-P} \leftrightarrow B_2^q, \quad O_3^q = (\bar{b}_i q_j)_{S-P} (\bar{b}_j q_i)_{S-P} \leftrightarrow B_3^q,$$

$$\{O_1^q, O_2^q\} \rightarrow \{O_1^q, O_3^q\}$$

Old basis \rightarrow New basis

Using the NLO and $O(1/m_b)$ relation:

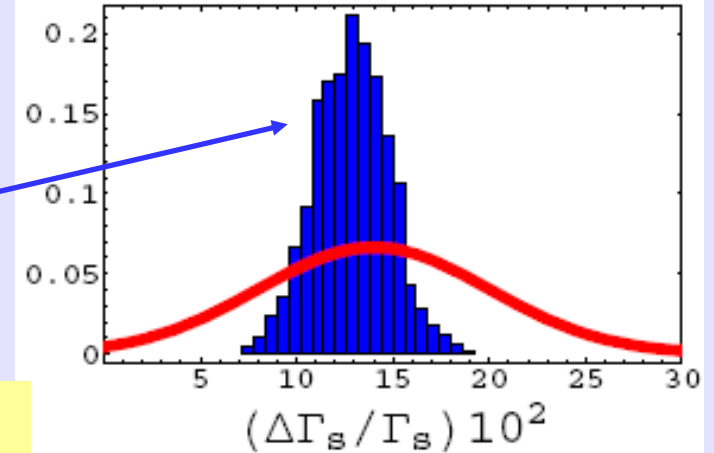
$$O_2^q + \alpha_1 O_3^q + \frac{\alpha_2}{2} O_1^q = R_0^q$$

	$O(1)$	$O(\alpha_s)$	$O(1)$	$O(\alpha_s)$	$O(1/m_b)$
$\text{Re} \left(\frac{\Gamma_{12}^s}{M_{12}^s} \right)^{\text{OLD}}$	\downarrow	\downarrow	\downarrow	\downarrow	\downarrow
	$= (-0.54 + 0.26) B_1^s$	$+ (-12.0 + 4.4) B_2^s$	$+ 4.0$		
$\text{Re} \left(\frac{\Gamma_{12}^s}{M_{12}^s} \right)^{\text{NEW}}$	\downarrow	\downarrow	\downarrow	\downarrow	\downarrow
	$= (-6.8 + 1.6) B_1^s$	$+ (-2.4 + 0.2) B_3^s$	$+ 1.6$		

The series in both
 α_s and $1/m_b$
 behave better!

	Old basis	New basis
$\frac{\Delta\Gamma_d}{\Gamma_d}$	$2.3(8)10^{-3}$	$4.1(5)10^{-3}$
$\frac{\Delta\Gamma_s}{\Gamma_s}$	$7(3)10^{-2}$	$13(2)10^{-2}$

NEW th. vs exp.



**Lenz and Nierste
quote $(12.7 \pm 2.4) \cdot 10^{-2}$:
perfect compatibility!**

**The non-perturbative uncertainties
have a smaller impact in the new basis**

BUT

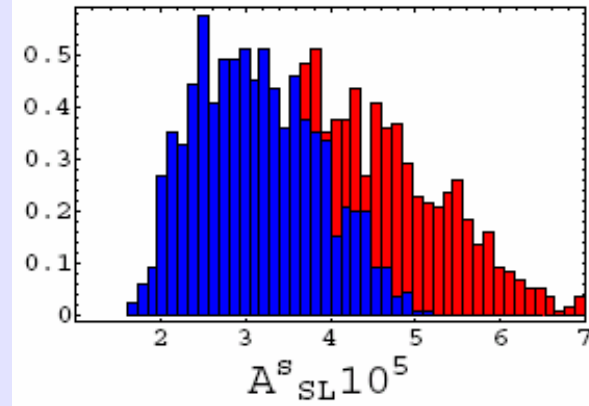
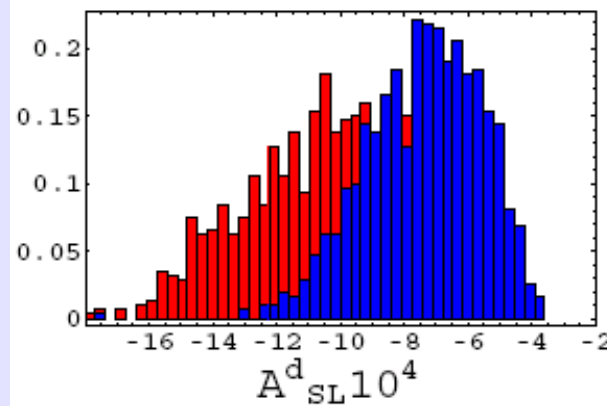
**The shift of central values signal important
(unknown) $O(\alpha_s^2)$ and $O(\alpha_s) \cdot O(1/m_b)$ corrections**

WHILE

**$O(1/m_b^5)$ effects are small
[A. Lenz and U. Nierste:
private communication]**

Semileptonic CP-Asymmetries

NLO distr.
VS
LO distr.



Experimental averages:

$$A_{\text{SL}}^{\text{d}} = -(30 \pm 78) \cdot 10^{-4}, \quad A_{\text{SL}}^{\text{s}} = (2450 \pm 1930(\text{stat.}) \pm 350(\text{syst.})) \cdot 10^{-5}$$

[HFAG, 2006]

[D0, 2006]

Theoretical predictions at the NLO + contribution of $\mathcal{O}(1/m_b^4)$:

$$A_{\text{SL}}^{\text{d}} = -(6.4 \pm 1.6) \cdot 10^{-4}, \quad A_{\text{SL}}^{\text{s}} = (2.7 \pm 0.6) \cdot 10^{-5}$$

[M.Ciuchini, E.Franco, V.Lubicz, F.Mescia, C.T., 2003]

In the combinations of
Wilson coefficients entering here,
the change of basis doesn't help

	Old basis	New basis
A_{SL}^{d}	$-6.4(16)10^{-4}$	$-6.6(17)10^{-4}$
A_{SL}^{s}	$2.7(6)10^{-5}$	$2.8(6)10^{-5}$

Outlook

We are waiting for updated experimental values of:

$$\frac{\tau(\Lambda_b)}{\tau(B_d)}$$

$$\frac{\Delta\Gamma_s}{\Gamma_s}$$

...and measurements with ~30% accuracy of:

$$\frac{\Delta\Gamma_d}{\Gamma_d}$$

$$A_{SL}^d$$

$$A_{SL}^s$$

2007

Meanwhile, from the theoretical side:

- Results for the perturbative $O(1/m_b^5)$ contribution to Γ_{12}^q [A.Lenz, U.Nierste]
- Results in QCD-sum rules for $\Delta B=2$ subleading operators [A.A.Pivarov et al.]
- Further Lattice studies of $\Delta B=2$ B-parameters [D.Becirevic et al,...]

...and hopefully soon:

- Improved lattice calculation of Λ_b B-parameters
- First Lattice calculation of $\Delta B=2$ subleading operators (R_2^q, R_3^q)
- Perturbative calculation of $O(\alpha_s^2)$ and $O(\alpha_s)O(1/m_b)$ corrections in Γ_{12}^q

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