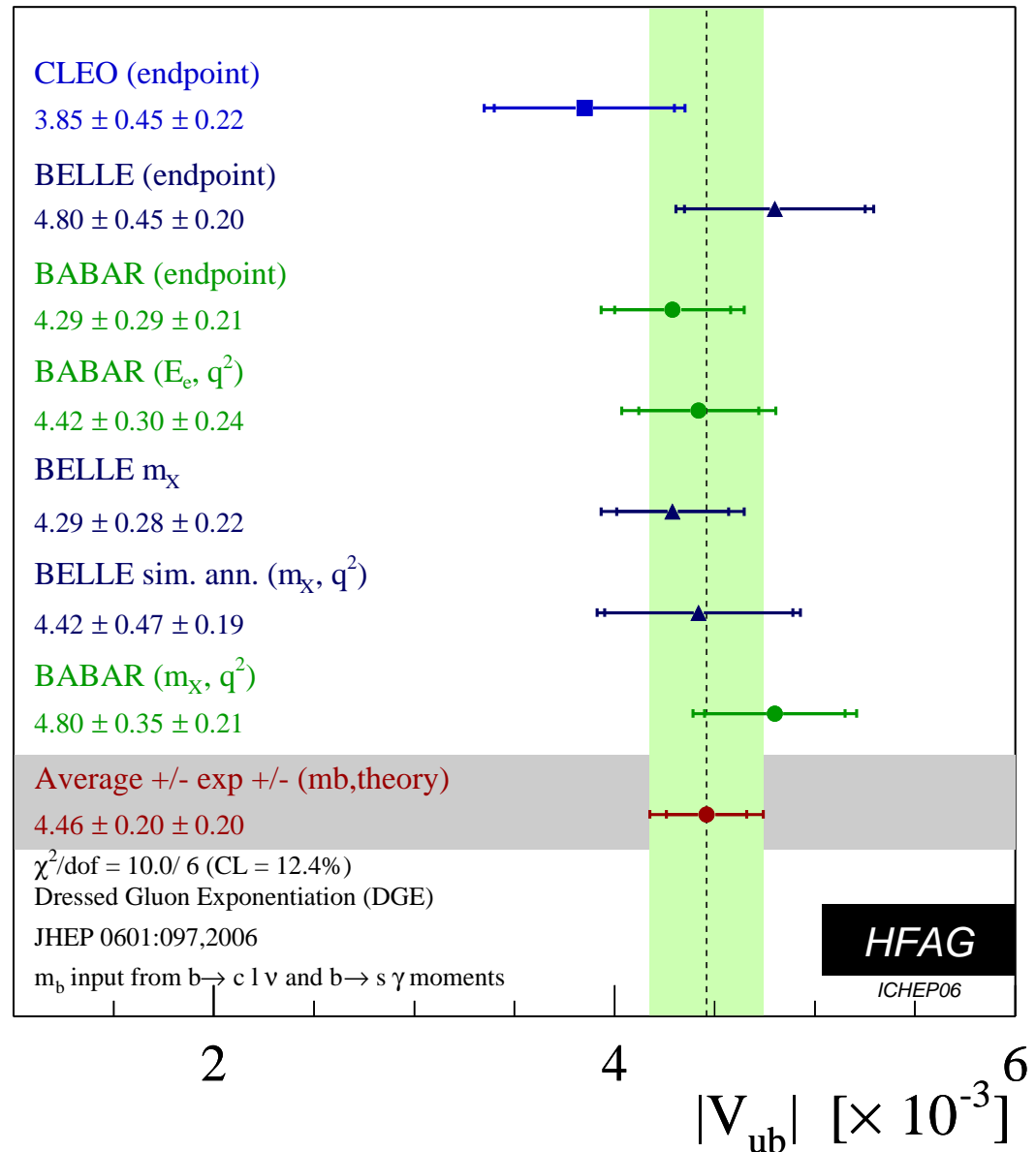


Einan Gardi (Cambridge)

- Inclusive $\bar{B} \rightarrow X_u l \bar{\nu}$ spectra
 - The essential rôle of resummation
 - Dressed Gluon Exponentiation (DGE) vs. shape function
- Recent progress and prospects for inclusive $\bar{B} \rightarrow X_u l \bar{\nu}$ spectra
 - Power corrections in the DGE approach
implemented in the new release
 - NNLO corrections
The $\mathcal{O}(\beta_0 \alpha_s^2)$ corrections to the fully differential width have been computed
Gambino, Gardi & Ridolfi [hep-ph/0610140]
- Lessons from $\bar{B} \rightarrow X_s \gamma$ DGE@NNLO analysis
Andersen & Gardi [hep-ph/0609250]

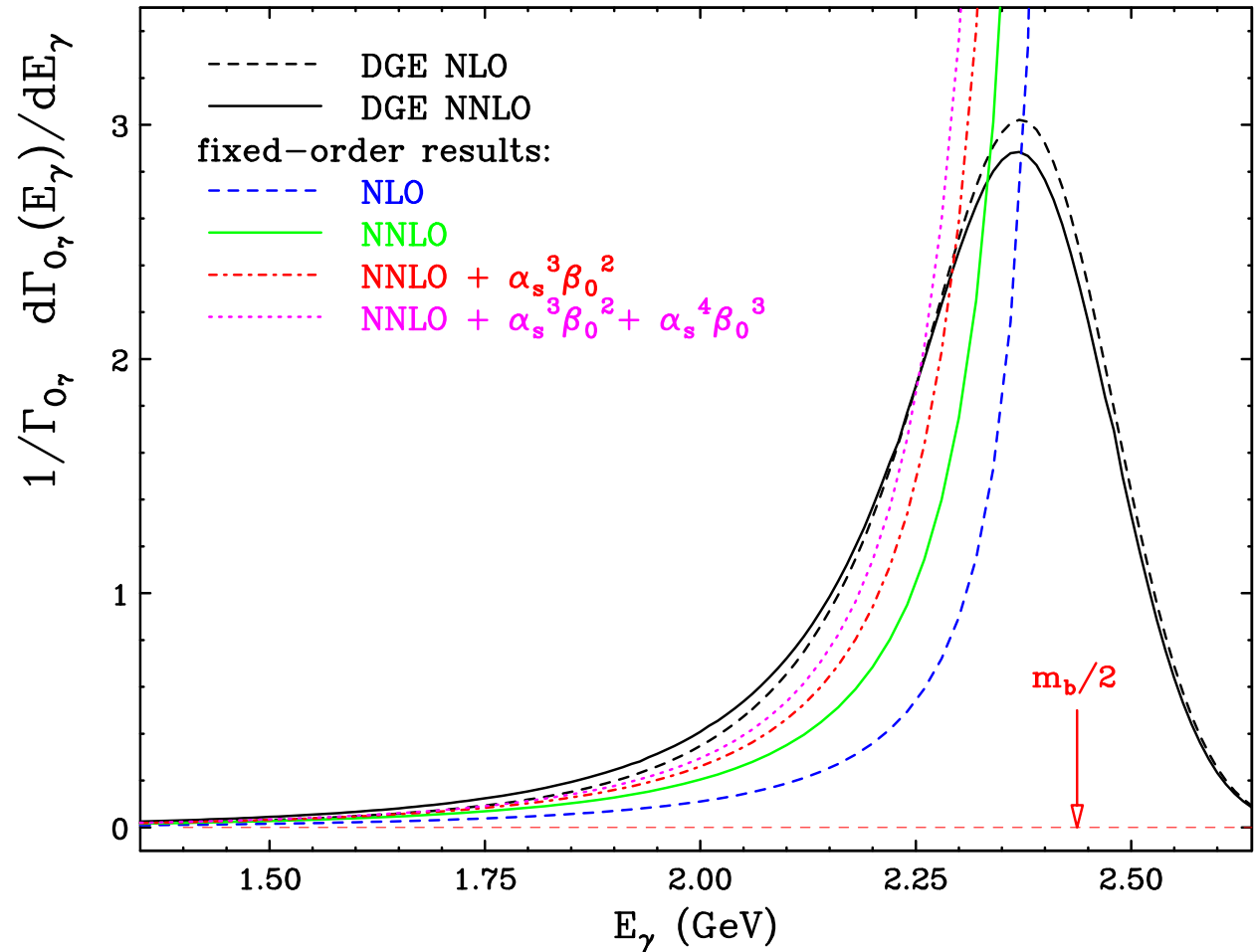
World Average $|V_{ub}|$ using DGE

- CLEO, Belle & BaBar have performed several inclusive measurements of the partial $b \rightarrow u$ BF with different kinematic cuts on E_l , q^2 , M_X , etc.
- **March '06**: each measurement was translated by *HFAG* into a value for $|V_{ub}|$ using **DGE**.
HFAG update: **July '06**
- The results are all consistent.
- The uncertainty is the smallest obtained so far.



$\bar{B} \rightarrow X_s \gamma$: DGE vs. the fixed-order expansion

- Stability
- Support



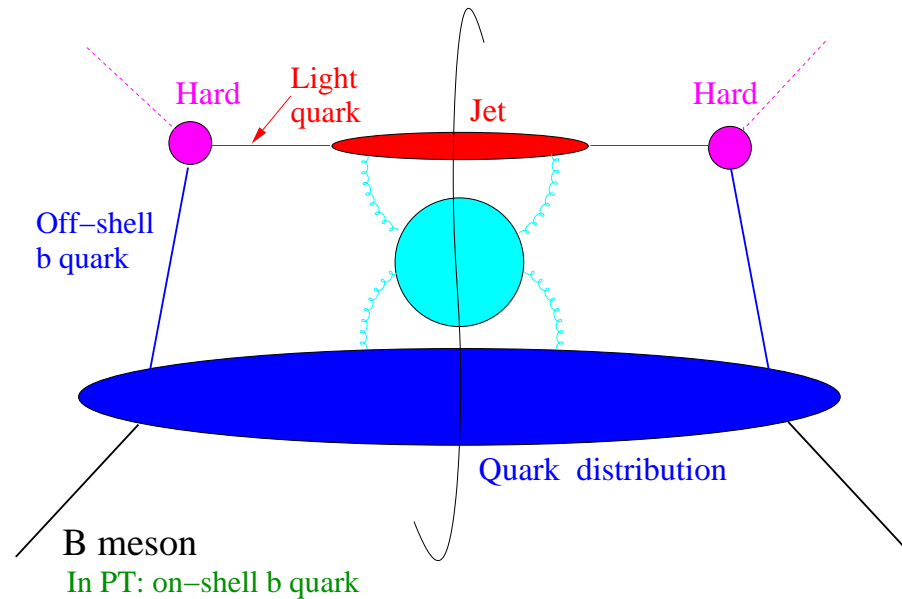
Inclusive $\bar{B} \rightarrow X_u l \bar{\nu}$ spectra by Dressed Gluon Exponentiation

- Inclusive decays spectra are **IR & Collinear safe**:
 - Infrared cutoff and leading–power shape function are not necessary.
 - All non-perturbative effects can be taken into account as power corrections in moment space.
- Large perturbative corrections are associate with
 - Sudakov logarithms
 - running–coupling effects (renormalons)

Both are resummed by DGE.
- General relation between the divergence of the perturbative expansion (**renormalons**) and infrared sensitivity (**power corrections**).
- In inclusive decays: resummation reveals that **IR & Collinear safety extends beyond the logarithmic level!**
 - Non-perturbative effects of the quark distribution first appear at $\mathcal{O}((N\Lambda/m_b)^3)$.
 - DGE spectrum **without any power corrections** is already a good approximation.
 - Given the known parametric dependence, power corrections can be effectively parametrized and quantified.

Factorization and DGE

Factorization in inclusive decays:
 Korchemsky & Sterman '94



DGE:

Dressed Gluon Exponentiation:
 (moment space)

<u>Hard:</u>		<u>Jet:</u>		<u>Quark Distribution:</u>
m_b	\gg	m_b/\sqrt{N}	\gg	m_b/N
calculated		calculated		calculated+p.c.**

** power corrections distinguishing $S(m_b/N; \mu)$ (the quark distribution in the meson) from $S_{\text{PT}}(m_b/N; \mu)$ (the quark distribution in an on-shell quark) are parametrized.

DGE vs. shape function

DGE:

Dressed Gluon Exponentiation: (moment space)	<u>Hard:</u> m_b calculated	\gg	<u>Jet:</u> m_b/\sqrt{N} calculated	\gg	<u>Quark Distribution:</u> m_b/N calculated+p.c.**
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** power corrections distinguishing $S(m_b/N; \mu)$ (the quark distribution in the meson) from $S_{\text{PT}}(m_b/N; \mu)$ (the quark distribution in an on-shell quark) are parametrized.

Cutoff-based “Shape Function” approach (BLNP):

“Shape Function” approach: (momentum space)	<u>Hard:</u> m_b calculated	\gg	<u>Jet:</u> $\sqrt{m_b\Lambda}$ calculated	\gg	<u>Quark Distribution:</u> Λ parametrized*
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* the entire leading-power quark distribution in the meson is parameterized.

Power corrections in the DGE approach

- Resummed perturbation theory (**on-shell** heavy quark) yields:

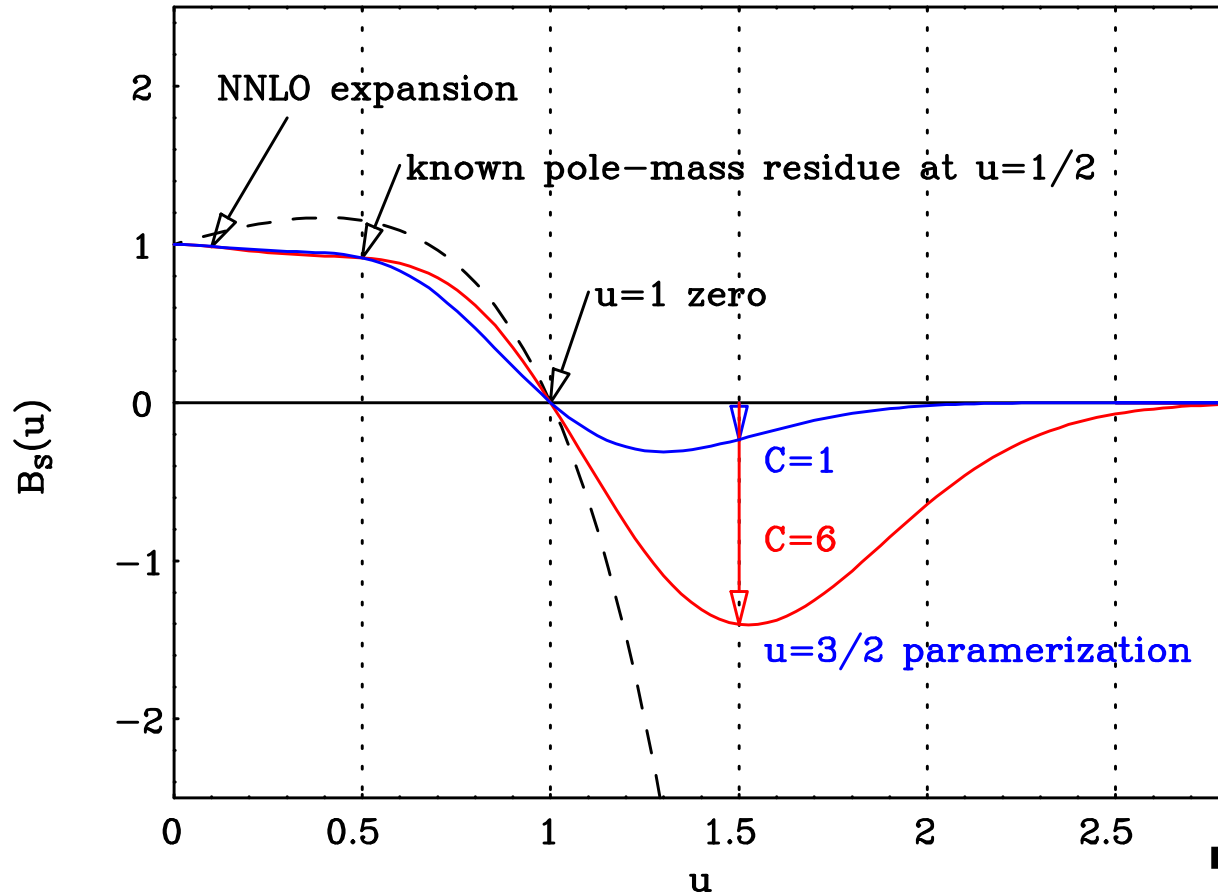
$$\text{Sud}(p^-, N) = \exp \left\{ \frac{C_F}{\beta_0} \int_0^\infty \frac{du}{u} \left(\frac{\Lambda}{p^-} \right)^{2u} \left[\underbrace{B_{\mathcal{S}}(u) \left(\frac{\Gamma(N)\Gamma(-2u)}{\Gamma(N-2u)} + \frac{1}{2u} \right)}_{\text{Quark Distribution: } N^{2u}} - \underbrace{B_{\mathcal{J}}(u) \left(\frac{\Gamma(N)\Gamma(-u)}{\Gamma(N-u)} + \frac{1}{u} \right)}_{\text{Jet: } N^u} \right] \right\}$$

- Renormalon ambiguities indicate power corrections $(N\Lambda/p^-)^k$ in the exponent!
 - $u = 1/2$ renormalon ambiguity cancels with the pole mass renormalon.
 - $u = 1$ renormalon is missing ($B_{\mathcal{S}}(1) = 0$).
 - $u \geq 3/2$ ambiguities are present in the on-shell decay spectrum.
- To approximate the quark distribution in the **meson**: introduce one non-perturbative parameter f_{PV} to parametrize $(N\Lambda/p^-)^k$ corrections for $k \geq 3$:

$$\text{Sud}(p^-, N) \longrightarrow \text{Sud}(p^-, N)|_{\text{PV}} \times \exp \left\{ \frac{C_F}{\beta_0} \pi f_{\text{PV}} \sum_{k=3}^{\infty} \frac{(-1)^k}{k k!} B_{\mathcal{S}}(k/2) \left(\frac{\Lambda}{p^-} \right)^k \prod_{j=1}^k (N-j) \right\},$$

The Sudakov factor (beyond fixed logarithmic accuracy)

$$\text{Sud}(p^-, N) \simeq \exp \left\{ \frac{C_F}{\beta_0} \int_0^\infty \frac{du}{u} \left(\frac{\Lambda}{p^-} \right)^{2u} \left[B_{\mathcal{S}}(u) \Gamma(-2u) (N^{2u} - 1) - B_{\mathcal{J}}(u) \Gamma(-u) (N^u - 1) \right] \right\}$$



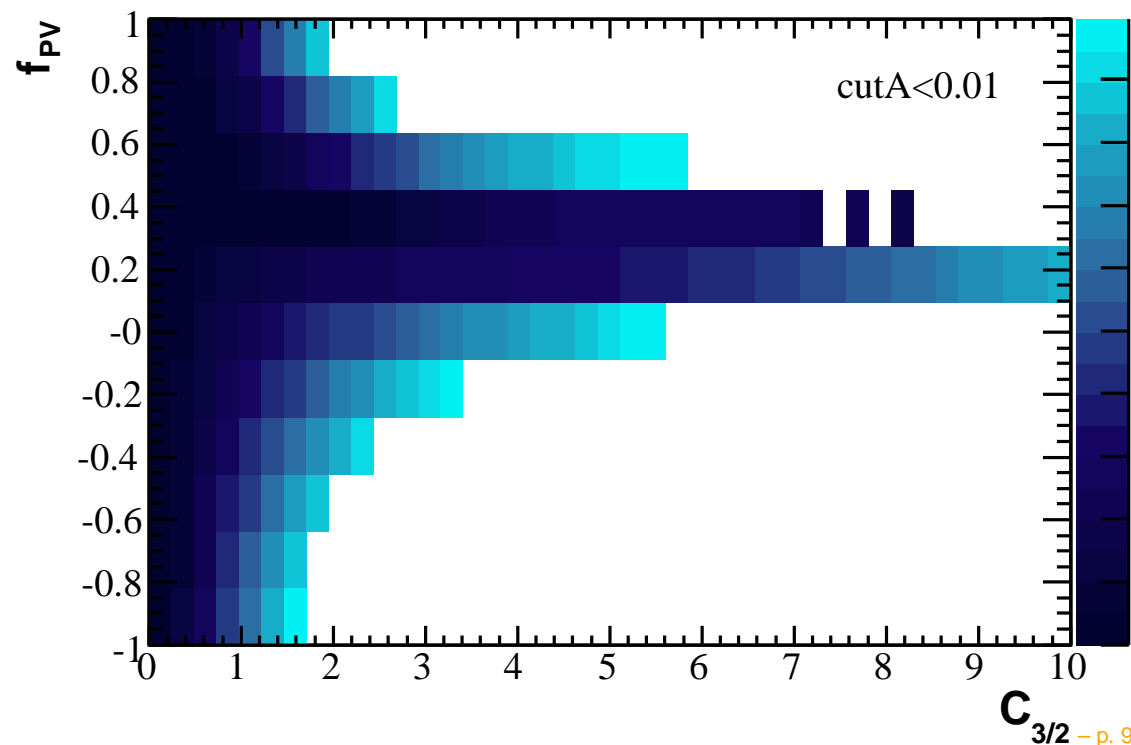
Power corrections in the DGE approach

- Simultaneous constraints on $B_S(u)$ (the renormalon residue $C_{3/2}$) and f_{PV} based on
 - support (done)
 - data (fits to $\bar{B} \rightarrow X_s \gamma$ moments under way)
- $\bar{B} \rightarrow X_s \gamma$ values of $C_{3/2}$ and f_{PV} directly apply to $\bar{B} \rightarrow X_u l \bar{\nu}$.
- Variation of $R_{cut} = \text{BF}(M_X < 1.7 \text{ GeV}) / \text{BF}(\text{total})$ due to power corrections: $\pm 1\%$

Support-based exclusion map

Francesca Di-Lodovico (BaBar)

Andersen & Gardi [hep-ph/0609250]



NNLO corrections

- Lesson from $\bar{B} \rightarrow X_s \gamma$ DGE@NNLO analysis:
 - significant reduction of uncertainty along the tail upon including $\mathcal{O}(\alpha_s^2)$ corrections (was done for O_7).
 - renormalization scale in the matching coefficient in $\bar{B} \rightarrow X_s \gamma$: $\mu \sim m_b/2$ — stability NLO \rightarrow NNLO
- $\beta_0 \alpha_s^2$ corrections dominate the $\mathcal{O}(\alpha_s^2)$ contribution to inclusive decays (both within the Sudakov region and out of it).
- The $\beta_0 \alpha_s^2$ corrections (as well as $\beta_0^n \alpha_s^{n+1}$ for any n) are now **available** for the fully differential $\bar{B} \rightarrow X_u l \bar{\nu}$ width.
Gambino, Gardi & Ridolfi [hep-ph/0610140].
- Matching to $\mathcal{O}(\beta_0 \alpha_s^2)$ will soon be included in the DGE code.
- Estimated variation of $R_{\text{cut}} = \text{BF}(M_X < 1.7 \text{ GeV}) / \text{BF}(\text{total})$ due to NNLO corrections: $\pm 6\%$

Conclusions

DGE provides a systematic and effective framework to compute B-meson decay spectra.

Status of DGE-based experimental studies:

- $|V_{ub}|$ has been determined from all available $\bar{B} \rightarrow X_u l \bar{\nu}$ BF measurements with a variety of cuts: consistent values; small error.
- Agrees well with $\bar{B} \rightarrow X_s \gamma$ moments as a function of the cut;
Fits to $\bar{B} \rightarrow X_u l \bar{\nu}$ moments or spectrum have not yet been done (Weak Annihilation).

Recent Progress:

- New way to include and constrain(!) power corrections
 - New release of the DGE code is available for both $\bar{B} \rightarrow X_s \gamma$ and $\bar{B} \rightarrow X_u l \bar{\nu}$.
 - Fits to data are necessary to determine m_b , $C_{3/2}$ and f_{PV} .
- Available NNLO $\beta_0 \alpha_s^2$ corrections for the fully differential $\bar{B} \rightarrow X_u l \bar{\nu}$ width.

Prospects:

- Matching to $\mathcal{O}(\beta_0 \alpha_s^2)$: reduction of theoretical uncertainty on $|V_{ub}|$ from any experimental cut.