Lattice QCD results for the $B \rightarrow D^{(*)} | v$ form factors: F(1) and G(1)

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Why study semileptonic $B \rightarrow D^{(*)}$ decays?



$$\frac{d\Gamma(B \to Dl\nu)}{dw} = \frac{G_F^2}{48\pi^3} m_D^3 (m_B + m_D)^2 (w^2 - 1)^{3/2} |V_{cb}|^2 |\mathcal{F}_{B \to D}(w)|^2
\frac{d\Gamma(B \to D^* l\nu)}{dw} = \frac{G_F^2}{4\pi^3} m_{D^*}^3 (m_B - m_{D^*})^2 \sqrt{w^2 - 1} \mathcal{G}(w) |V_{cb}|^2 |\mathcal{F}_{B \to D^*}(w)|^2$$

$$w \equiv v' \cdot v
w = 1
at zero recoil$$

- ← Experiments can only measure the product (form factor) x |V_{cb}|
- Lattice QCD calculations needed to determine normalization and extract the CKM matrix element |V_{cb}|
- Only need one q² point from lattice -- choose w=1 because easiest to calculate

Lattice calculations of F(1), G(1)

QUENCHED

- Calculated by Fermilab collaboration [Phys.Rev.D66:014503,2002; Phys.Rev.D61:014502,2000]
 - Wilson u,d,s quarks; Fermilab formulation for b,c quarks

UNQUENCHED

- Calculations in progress by the Fermilab/MILC collaboration
- Use the publicly available "2+1 flavor" MILC configurations [Phys.Rev.D70:114501,2004] which have three flavors of improved staggered quarks:
 - * Two degenerate light quarks and one heavy quark ($\approx m_s$)
 - Light quark mass ranges from $m_s/10 \le m_l \le m_s$
- Fermilab formulation for b,c quarks
- ◆ Preliminary B→D form factor, G(1), presented at Lattice '04 [Nucl.Phys.Proc.Suppl.140:461-463,2005]

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Schematic of $B \rightarrow D$ form factor calculation



G(1) depends upon two form factors, h₊ and h_{_:}

$$\mathcal{F}_{B\to D}(1) = h_{+}^{B\to D}(1) - \frac{m_B - m_D}{m_B + m_D} h_{-}^{B\to D}(1)$$

- Could calculate directly from lattice B→D correlation function, but errors would be large
- Instead construct double ratios of correlation functions to cancel out bulk of statistical fluctuations from Monte Carlo simulations as well as many systematic errors

$$\frac{\langle D|\overline{c}\gamma^4 b|\overline{B}\rangle\langle\overline{B}|\overline{b}\gamma^4 c|D\rangle}{\langle D|\overline{c}\gamma^4 c|D\rangle\langle\overline{B}|\overline{b}\gamma^4 b|\overline{B}\rangle} = |h_+(1)|^2 \qquad \frac{\langle D|\overline{c}\gamma_i b|\overline{B}\rangle\langle D|\overline{c}\gamma_0 c|D\rangle}{\langle D|\overline{c}\gamma_0 b|\overline{B}\rangle\langle D|\overline{c}\gamma_i c|D\rangle} = 1 - \frac{h_-(w)}{h_+(w)}$$

◆ Can calculate h+ and h. directly at the tuned bottom & charm quark masses

Schematic of $B \rightarrow D^*$ form factor calculation



+ F(1) depends upon one form factor, h_{A_1} :

 $\mathcal{F}_{B\to D^*}(1) = h_{A_1}(1)$

- ♦ h_{A1}, however, cannot be determined from a single double ratio
- Instead use heavy quark effective theory (HQET) as an intermediary ...

Schematic of $B \rightarrow D^*$ form factor calculation

- Can expand $h_{A_1}(1)$ in powers of the heavy quark mass using HQET
- To order 1/m_Q², depends upon three different HQET matrix elements:

$$h_{A_1}(1) = \eta_A \left[1 - \frac{\ell_V}{(2m_c)^2} + \frac{2\ell_A}{2m_c m_b} - \frac{\ell_P}{(2m_b)^2} \right]$$

- Can determine three HQET matrix elements from the heavy quark mass dependence of three different double ratios
- For example, use the heavy-quark mass dependence of R_+ to extract I_p :

$$\mathcal{R}_{+} = \frac{\langle D | \overline{c} \gamma^{4} b | \overline{B} \rangle \langle \overline{B} | \overline{b} \gamma^{4} c | D \rangle}{\langle D | \overline{c} \gamma^{4} c | D \rangle \langle \overline{B} | \overline{b} \gamma^{4} b | \overline{B} \rangle} = |h_{+}(1)|^{2} \qquad h_{+}(1) = \eta_{V} \left[1 - \ell_{P} \left(\frac{1}{2m_{c}} - \frac{1}{2m_{b}} \right)^{2} \right]$$

+ Then recombine the HQET matrix elements get the form factor $h_{A_1}(1)$

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Comparison of F(1), G(1) lattice calculations



- Essentially same code to calculate correlation functions
- Both calculations use double ratios to reduce errors

DIFFERENCES

- ♦ G(1) comes directly from lattice data at tuned bottom & charm masses
- F(1) requires intermediate step of determining HQET matrix elements, and consequently lattice data at multiple heavy quark masses

Systematics in lattice calculations

- Lattice calculations typically quote the following sources of error:
 - 1. Monte carlo statistics & fitting
 - 2. Tuning lattice spacing, a, and quark masses
 - 3. Matching lattice gauge theory to continuum QCD
 - 4. Extrapolation to continuum

5. Chiral extrapolation to physical up, down quark masses

 Some lattice simulations also neglect dynamical quark loops -- known as the "quenched approximation"

 Errors #3 and #5 are dominant sources of uncertainty in current heavy-light lattice calculations -- will discuss them in turn

Heavy quarks on the lattice

PROBLEM: Generic lattice quark action will have discretization errors $\propto (am_Q)^n$

SOLUTION: Fermilab method uses knowledge of the heavy quark limit of QCD to systematically eliminate HQ discretization errors order-by-order [Phys.Rev.D55:3933-3957,1997; Phys.Rev.D62:014505,2000; Phys.Rev.D65:094514,2002]



- Requires tuning parameters of lattice action and lattice currents to continuum
- Typically calculate matching coefficients in lattice perturbation theory [Phys.Rev.D48:2250-2264,1993]
- Combine all errors associated with discretizing action into "heavy quark discretization errors"
- Estimate errors using knowledge of short-distance coefficients and powercounting

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Chiral extrapolation of lattice data

- Must extrapolate lattice results to physical values of up, down quark mass
- For MILC 2+1 flavor lattices, must use staggered chiral perturbation theory [Lee & Sharpe, Aubin & Bernard, Sharpe & RV]
 - * Accounts for next-to-leading order light quark mass dependence and light quark discretization effects through $O(\alpha_s^2 a^2 \Lambda_{QCD}^2)$
- For B→D, extrapolation under good
 control [Phys.Rev.D73:054501,2006]



black = finite lattice spacing red = continuum extrapolation For B→D^{*}, extrapolation using correct
 xPT is essential [Laiho, Lattice '06]



cusp at $m_{\pi^2} = (m_{D^*} - m_D)^2$

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Quenched results for F(1), G(1)

Quenched results for $B \rightarrow D^{(*)}$ form factors



- 4% determination of F(1) [Phys.Rev.D66:014503,2002]
- 2% determination of G(1) [Phys.Rev.D61:014502,2000]
- ◆ Proof-of-method: use of double ratios allows precise calculation of B→D^(*) form factors at zero recoil

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Potential sources of improvement

NECESSARY:

- Unquenched calculation with three light quark flavors (in progress)
- Use chiral perturbation theory expressions for extrapolation to physical up/ down quark mass (Staggered chiral perturbation theory for F(1) and G(1) calculated by Laiho & RV)

HELPFUL:

- Increased statistics (in unquenched calculation)
- + Lighter quark masses (in unquenched calculation) and finer lattice spacings
- ◆ 2-loop perturbative (or nonperturbative) matching
- Improved heavy-quark action (in progress -- Kronfeld & Oktay)
- Heavy-light calculations with different light quark action, e.g domain-wall (RBC) or overlap fermions (JLQCD)

Progress in unquenched calculations of F(1) and G(1)

Unquenched result for $B \rightarrow D$ form factor

$$\mathcal{F}_{B\to D}^{n_f=2+1}(1) = 1.074 \ (18)_{\text{sta}}(16)_{\text{sys}}$$
$$|V_{cb}| = 3.8 \ (1)_{\text{sta}}(6)_{\text{sys}} \times 10^{-2}$$

[Nucl.Phys.Proc.Suppl,140:461-463,2005] (preliminary)

Consistent with earlier quenched result

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- Linear extrapolation in quark mass and single lattice spacing
- G(1) calculation in progress with many improvements:
 - 4 times statistics

-- still only 2% error

- Additional data point at even lighter quark mass
- Extrapolation with staggered chiral perturbation theory expression
- Additional (larger) lattice spacing to estimate lattice spacing dependence

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"Prediction" for errors in $B \rightarrow D^*$ form factor

- Total error in quenched F(1) is 4%
- In addition to removing quenching uncertainty, unquenched calculation will make the following improvements:
 - Increased statistics
 - Extrapolation in light quark mass with staggered chiral perturbation theory
 - Because staggered χPT includes light quark discretization effects, will also reduce lattice spacing error
- Error in unquenched F(1) will likely be 2-3%
- To get below this, will help to reduce experimental uncertainty in the D-D^{*}-π coupling, g_π
- Size of continuum cusp varies
 by ± 1% when g_π is varied within its experimental uncertainty



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Summary and prospects for exclusive |V_{cb}|

- ♦ |V_{cb}| currently known to 4% from exclusive decays [PDG]
- ← Limited by theoretical uncertainty in the $B \rightarrow D^{(*)}$ form factors at zero recoil
- Quenched calculations demonstrate the capability of lattice calculations to determine F(1), G(1) to few percent accuracy
- Unquenched calculations are in progress using the same methodology
- Although the errors will likely only go from 4% (+ ???%)→2-3% between the quenched and unquenched calculations, the results will be on a stronger theoretical footing, and hence more reliable
- ✦ Reducing the errors to below 2-3% will require additional work, e.g.
 - Higher-order matching
 - * Better determination of g_{π}
- Unquenched calculations are happening right now, so stay tuned...