

# Enhanced Nonlocal Power Corrections to the $B \rightarrow X_s \gamma$ Decay Rate

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# Introduction

- Precision studies of inclusive  $B \rightarrow X_s \gamma$  decay are a cornerstone of quark-flavor physics
- Theoretical analysis of total inclusive rates based on local operator product expansion, by which total rates are expressed in terms of local HQET matrix elements; power corrections enter first at order  $(\Lambda_{QCD}/m_b)^2$
- Leading term now computed (almost) at NNLO in RG-improved perturbation theory [→ talks by P. Gambino, T. Becher]



# Introduction

- OPE breaks down for differential decay distributions near phase-space boundaries
- Twist expansion involving forward Bmeson matrix elements of nonlocal operators (shape functions) required to systematically account for nonperturbative effects [→ talk by G. Paz]



# Introduction

- Well-known that OPE faces limitations in B→X<sub>s</sub>γ decay, e.g.:
  - Charm-penguin contribution leads to  $\lambda_2/m_c^2$  correction when charm is treated as a heavy quark [Voloshin]



– Effect described by a nonlocal subleading shape function when counting  $m_c{}^2 \gg m_b \Lambda_{\rm QCD}$  is adopted



- Power corrections to high-energy part of B→X<sub>s</sub>γ photon spectrum can be parameterized systematically in terms of subleading shape functions
- Some of them (those considered so far) reduce to local operators when photon spectrum is integrated over energy
- *But:* Several others do not! [Lee, MN, Paz (in prep.)]



• Example of contributions to discontinuity of hadronic tensor:





 First two graphs give rise to 4-quark operators containing light u,d quarks besides s quarks



 $\rightarrow$  expected to give dominant effect for B<sup>0</sup>, B<sup>+</sup>

 Due to interference with Q<sub>7γ</sub>, effect is centered at large photon energy (cannot be eliminated by a cut)



• Contribution to total rate parameterized by matrix elements of tri-local operators:

$$\Delta \Gamma = -\Gamma_{77} \frac{C_{8g}}{C_{7\gamma}} \frac{4\pi \alpha_s}{N_c m_b} \int_{-\infty}^0 ds \int_{-\infty}^0 dt$$
$$\times \langle \bar{B} | C_F \left( O_1 + O_2 \right) - \left( T_1 + T_2 \right) | \bar{B} \rangle$$

Where:  $O_{1} = \sum_{q} e_{q} \bar{h}_{v}(0) \Gamma_{R} q(t\bar{n}) \bar{q}(s\bar{n}) \Gamma_{R} h_{v}(0),$  $O_{2} = \sum_{q} \frac{e_{q}}{2} \bar{h}_{v}(0) \Gamma_{R} \gamma_{\perp \alpha} q(t\bar{n}) \bar{q}(s\bar{n}) \gamma_{\perp}^{\alpha} \Gamma_{R} h_{v}(0),$  $T_{1} = \sum_{q} e_{q} \bar{h}_{v}(0) \Gamma_{R} t_{a} q(t\bar{n}) \bar{q}(s\bar{n}) \Gamma_{R} t_{a} h_{v}(0),$  $T_{2} = \sum_{q} \frac{e_{q}}{2} \bar{h}_{v}(0) \Gamma_{R} \gamma_{\perp \alpha} t_{a} q(t\bar{n}) \bar{q}(s\bar{n}) \gamma_{\perp}^{\alpha} \Gamma_{R} t_{a} h_{v}(0)$ 

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# Relation with Optical Theorem

- Normally, expect 4-quark operators to contribute at order (Λ<sub>QCD</sub>/m<sub>b</sub>)<sup>3</sup> to the total decay rate - why different here?
- Reason is that in  $B \rightarrow X_s \gamma$  one does not sum over all cuts:



• Process is not really inclusive (partonic substructure of photon)



- Reliable field-theoretic estimates of these effects are very difficult to obtain (lattice QCD unable to handle light-cone operators)
- Naïve dimensional analysis suggests:  $\Delta\Gamma/\Gamma_{77} \sim (C_{8g}/C_{7\gamma}) \pi \alpha_s (\Lambda/m_b)$
- Could be a 5% correction to the rate!



 Our approach: Apply vacuum insertion approximation

 motivated by large-N<sub>c</sub> counting rules
 well tested for local 4-quark operators in analysis of B-hadron lifetimes
 not uncontroversial for nonlocal operators (not clear to me why ...)



- Matrix elements of operators O<sub>2</sub>,T<sub>1</sub>,T<sub>2</sub> vanish in VIA due to color and/or Dirac structure
- Matrix element of O<sub>1</sub> gets related to leading B-meson light-cone distribution amplitude in position space:

$$egin{aligned} &\langle ar{B} | \, O_1 \, | ar{B} 
angle_{ ext{VIA}} = e_q \, rac{f_B^2 m_B}{4} \, \widetilde{\phi}^B_+(s) \, [\widetilde{\phi}^B_+(t)]^* \end{aligned}$$



• Integration over light-cone gives inverse moment of LCDA (parameter  $\lambda_B$ ):

$$-i\int_{-\infty}^{0}ds\,\widetilde{\phi}^{B}_{+}(s)=\int_{0}^{\infty}rac{d\omega}{\omega}\,\phi^{B}_{+}(\omega)=rac{1}{\lambda_{B}}$$

• Result:

$$egin{aligned} &\Delta \Gamma_{ ext{VIA}} \ &\Gamma_{ extsf{77}} = -rac{e_q C_{8g}}{C_{7\gamma}} \, rac{\pi lpha_s}{2} \left(1 - rac{1}{N_c^2}
ight) rac{f_B^2 m_B}{\lambda_B^2 m_b} \ &pprox -0.26 e_q \left(rac{f_B}{\lambda_B}
ight)^2 \end{aligned}$$



- Effect potentially large
- Leading contribution to flavor-dependent rate asymmetry:

$$\frac{\Gamma(B^- \to X_s \gamma) - \Gamma(\bar{B}^0 \to X_s \gamma)}{\Gamma(\bar{B} \to X_s \gamma)} \approx -0.05 \left(\frac{\lambda_B}{0.5 \,\text{GeV}}\right)^{-2}$$

- Expect effect between -2% and -20% (sign determined)
- Contribution to flavor-averaged decay rate is 6 times smaller (few % effect at most in VIA)



#### Caveats!

- These estimates are very model dependent (even though reasonable)
- They do not provide conservative error estimates!
- We believe uncertainty on total  $B \rightarrow X_s \gamma$ branching ratio should not exceed 5%, but this is difficult to prove



## Conclusions

- There is no local OPE for the total  $B{\rightarrow}X_{s}\gamma$  decay rate!
- Nonlocal power corrections (subleading shape functions) starting at order  $\Lambda_{\rm QCD}/m_{\rm b}$  exist even for the total decay rate
- Irreducible theory error, of similar magnitude as perturbative uncertainty at NNLO
- Measurement of flavor-dependent asymmetry could help to corroborate our model estimates